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# On instant blow-up for quasilinear parabolic equations with growing initial data

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We are interested in the existence of the solutions of the parabolic equations with initial data which are not bounded at space infinity.

In [4] Giga and the author considered a nonnegative blowing up solution of the semilinear parabolic equation of the form

$$u_t = \Delta u + f(u), \quad x \in \mathbf{R}^N, t > 0$$

with nonlinear terms  $f$  and nonnegative initial data  $u_0$  satisfying that  $f$  is positive, nondecreasing and convex in  $(0, \infty)$ ,  $\int_1^\infty ds/f(s) < \infty$  and there are sequences  $\{x_n\} \subset \mathbf{R}^N$  and  $\{r_n\} \subset \mathbf{R}_+$  with  $\lim_{n \rightarrow \infty} |x_n| = \infty$  and  $\lim_{n \rightarrow \infty} r_n \geq 0$  such that

$$\lim_{n \rightarrow \infty} \frac{b_n}{r_n^2 f(b_n)} \text{ is small enough}$$

with  $b_n = \inf\{u_0(x) : |x - x_n| \leq r_n\}$ . They showed that the solutions do not exist even locally in time.

We consider the initial value problem for a quasilinear parabolic equation of the form

$$\begin{cases} u_t = \Delta u^m + u^p, & x \in \mathbf{R}^N, t \in (0, T), \\ u(x, 0) = u_0(x), & x \in \mathbf{R}^N. \end{cases} \quad (1)$$

Here we assume that  $N \geq 1$ ,  $1 \leq m < p$ .

We are interested in the problem whether there is a local-in-time solution of (1) when an initial datum  $u_0$  is continuous and grows at the space infinity, for example  $\lim_{|x| \rightarrow \infty} u_0(x) = \infty$ .

We consider the weak solution  $u$  in  $\mathbf{R}^N \times [0, T)$  of (1) such that  $u \in C(\mathbf{R}^n \times [0, \tau))$  for each  $\tau \in (0, T)$ , and for any bounded domain  $\Omega \in \mathbf{R}^N$  with smooth boundary  $\partial\Omega$ ,  $0 < \tau < T$  and nonnegative  $\phi(x, t) \in C^{2,1}(\Omega \times [0, T))$  which vanishes on the boundary  $\partial\Omega$ ,

$$\begin{aligned} & \int_{\Omega} u(x, \tau) \phi(x, \tau) dx - \int_{\Omega} u(x, 0) \phi(x, 0) dx \\ &= \int_0^{\tau} \int_{\Omega} \{u \partial_t \phi + u^m \Delta \phi + u^p \phi\} dx dt - \int_0^{\tau} \int_{\partial\Omega} u^m \partial_{\nu} \phi dS dt, \end{aligned} \quad (2)$$

where  $\nu$  denote the outer unit normal to the boundary. Note that the solution of (1) may be nonunique. Define  $T^* = T^*(u_0)$  as the supremum of all existence times of these solutions.

In this paper we shall prove that  $T^* = 0$  when the initial data  $u_0$  is growing at the space infinity. In other words there is even no local-in-time solution such that for any  $\tau > 0$  the weak solution does not exist for  $t \in (0, \tau)$ . We say this phenomenon  $T^* = 0$  an *instant blow-up*. We are able to prove that the instant blow-up occurs for more general initial data  $u_0$ .

**Theorem.** Assume that  $u_0 \in C(\mathbf{R}^N)$  is nonnegative. Assume that there are sequences  $\{x_n\}_{n=1}^{\infty} \subset \mathbf{R}^N$  and  $\{r_n\}_{n=1}^{\infty} \subset \mathbf{R}_+$  with  $\lim_{n \rightarrow \infty} |x_n| = \infty$  and  $\lim_{n \rightarrow \infty} r_n \geq 0$  such that

$$\lim_{n \rightarrow \infty} r_n^2 b_n^{p-m} > \frac{1}{\varepsilon} \quad (3)$$

for some  $\varepsilon \in (0, 1/c)$ , where  $b_n = \inf\{u_0(x) : |x - x_n| \leq r_n\}$  and  $c > 0$  is the first eigenvalue of  $-\Delta$  in a unit ball with the Dirichlet boundary condition. Then  $T^* = 0$ , i.e., the instant blow-up occurs provided that only nonnegative solutions are considered.

The proof of Theorem depends on a classical Kaplan's argument [6] to show the existence of blow-up which uses principal eigenfunctions of the Laplace operator with the Dirichlet condition.

In [1] among other results there is one about a sufficient condition on initial data for nonexistence of a local-in-time nonnegative solution for  $u_t = \Delta u^m + u^p/(1 + |x|)^{\alpha}$  with  $m \geq 1$ ,  $p > 1$  and  $\alpha \in \mathbf{R}$ . In the case of  $\alpha = 0$  the condition leads

$$\sup_{x \in \mathbf{R}^n} \int_{B(x,1)} u_0(y) dy = \infty. \quad (4)$$

In [1] this is explicitly mentioned for  $1 < p < m + 2/N$ . However, their proof is still valid for all  $p > 1$ . By the way their main interest is the existence of

solution; for example they proved the local existence when

$$\sup_{x \in \mathbf{R}^N} \int_{B(x,1)} u_0(y) dy < \infty$$

for  $1 < p < m + 2/N$ . The condition (3) is not included in the condition of their result for  $p > m + 2/N$ . In fact, if  $u_0 \geq b_n$  on  $B(x_n, r_n)$ , then  $\lim_{n \rightarrow \infty} b_n r_n^N = \infty$  is a sufficient condition for (4) (not a necessary condition). Our condition leads  $\lim_{n \rightarrow \infty} r_n^2 b_n^{p-m}$  is large enough. This shows that our condition for  $p > m + 2/N$  is not included in their condition.

In [1] they also prove the local existence for  $p \geq m + 2/N$  when  $u_0$  fulfills

$$\sup_{x \in \mathbf{R}^n} \int_{B(x,1)} u_0^q(y) dy < \infty$$

for some  $q > N(p - m)/2$ . In our nonexistence result  $u_0$  satisfies

$$\sup_{x \in \mathbf{R}^n} \int_{B(x,1)} u_0^q(y) dy \geq \lim_{n \rightarrow \infty} \int_{B(x_n,1)} u_0^q(y) dy \geq \lim_{n \rightarrow \infty} \varepsilon^{-\frac{q}{p-m}} r_n^{N-\frac{2q}{p-m}} = \infty$$

for any  $q > N(p - m)/2$ , where  $\varepsilon$  is used in (3).

In [4] Theorem was proved in the case  $m = 1$ . They studied the instant blow-up by using not only the eigenfunction method in [6] same as this paper but also the energy method in [7] and [2].

In the rest of the paper Theorem will be proved by using the Kaplan's argument [6].

**Lemma.** (c.f. [3, Lemma 4.2]) *Let  $v$  be the solution of the integral equation of the form*

$$v(t) - v(0) = \int_0^t h(v(s)) ds \quad (5)$$

*in  $[0, T_0)$  with  $h$  satisfying  $h \in C^1[0, \infty)$  and  $h' \geq 0$ . Let  $\tilde{v}$  be a nonnegative measurable function on  $[0, T_0)$ . Assume that  $\tilde{v}$  satisfies*

$$\tilde{v}(t) - \tilde{v}(t_0) \geq (\leq) \int_{t_0}^t h(v(s))(s) ds \quad \text{for } t_0, t \in [0, T_0) \quad \text{with } t_0 \leq t. \quad (6)$$

*Assume that  $\tilde{v}(0) \geq (\leq) v(0)$ . Then*

$$\tilde{v}(t) \geq (\leq) v(t) \quad \text{for } t \in [0, T_0).$$

*Proof.* We shall only prove the case  $\tilde{v}(t) - \tilde{v}(t_0) \geq \int_{t_0}^t \tilde{v}^p(s) ds$  since the proof of the other case is parallel. Since  $\tilde{v}(0) \geq v(0)$ , the estimate (6) together with (5) yields

$$\tilde{v}(t) - v(t) \geq \int_0^t (h(\tilde{v}(s)) - h(v(s))) ds.$$

By the mean value theorem we observe that

$$\tilde{v}(t) - v(t) \geq \int_0^t c(s) (\tilde{v}(s) - v(s)) ds,$$

where

$$c(s) = \int_0^1 h'(\theta v(s) + (1 - \theta)\tilde{v}(s)) d\theta.$$

We set  $\psi_\epsilon(t) = \tilde{v}(t) - v(t) + \epsilon$  with  $\epsilon > 0$ , and observe that  $\psi_\epsilon(t)$  satisfies

$$\psi_\epsilon \geq \int_0^t c(s) \psi_\epsilon(s) ds + \epsilon \left(1 - \int_0^t c(s) ds\right).$$

We set

$$t_1 = \sup \left\{ t > 0; \int_0^t c(s) ds < \frac{1}{2} \right\}.$$

Then, for  $t \in [0, t_1]$  we have

$$\psi_\epsilon(t) \geq \int_0^t c(s) \psi_\epsilon(s) ds + \frac{\epsilon}{2}. \quad (7)$$

We shall argue by contradiction to prove  $\psi_\epsilon(t) \geq 0$ . Suppose that  $\psi_\epsilon(t) < 0$  for some  $t \in [0, t_1]$ . Then  $\psi_\epsilon(\tau) = 0$  for

$$\tau = \inf \{t \in [0, t_1]; \psi_\epsilon < 0\}. \quad (8)$$

This  $\tau$  must be positive. Indeed, since  $\tilde{v}$  is nondecreasing by (6) and  $v$  is continuous,  $\psi_\epsilon(0) > \epsilon$  implies  $\tau > 0$ .

Since  $\int_0^\tau c(s) \psi_\epsilon(s) ds \geq 0$  and (8) imply  $\psi_\epsilon(\tau) \leq 0$ , we get a contradiction by (7). We thus proved that

$$\psi_\epsilon(t) \geq 0.$$

Since this holds for all  $\epsilon > 0$ , we get  $\tilde{v}(t) \geq v(t)$  for  $t \in [0, t_1]$ . (If  $\tilde{v}(t) < v(t)$  for some  $t$ , there exist  $\epsilon > 0$  such that  $\psi_\epsilon < 0$  for such  $t$ .)

Next, since  $\tilde{v}(t) \geq v(t)$  for  $t \in [0, t_1]$ , we observe that

$$\psi_\epsilon \geq \int_{t_1}^t c(s)\psi_\epsilon(s)ds + \epsilon \left(1 - \int_{t_1}^t c(s)ds\right).$$

We set

$$t_2 = \sup \left\{ t > t_0; \int_{t_1}^t c(s)ds < \frac{1}{2} \right\}$$

and observe that

$$\psi_\epsilon \geq \int_{t_1}^t c(s)\psi_\epsilon(s)ds + \frac{\epsilon}{2}$$

for  $t \in [t_1, t_2]$ . By the same argument one can prove  $\psi_\epsilon \geq 0$  for all  $\epsilon > 0$ , and  $\tilde{v}(t) \geq v(t)$  for  $t \in [t_1, t_2]$ .

We repeat this argument and conclude that

$$\tilde{v}(t) \geq v(t)$$

for all  $t \in [0, T_0)$ . By the same argument, we find if

$$\tilde{v}(t) - \tilde{v}(t_0) \leq \int_{t_0}^t \tilde{v}^p(s)ds \quad \text{for } t_0, t \in [0, T_0) \quad \text{with } t_0 \leq t,$$

then

$$\tilde{v}(t) \leq v(t) \quad \text{for } t \in [0, T_0).$$

□

*Proof of Theorem.* Let  $\{r_n\}_{n=1}^\infty$ ,  $\{x_n\}_{n=1}^\infty$  and  $\{b_n\}_{n=1}^\infty$  be as in Theorem satisfying (3). Set  $\lambda_n > 0$  denote the principal eigenvalue of  $-\Delta$  with Dirichlet problem in  $B(0, r_n)$ , and let  $\phi_n(x) \geq 0$  denote the corresponding positive eigenfunction normalized by  $\int_{B(0, r_n)} \phi_n(x)dx = 1$ . By scaling it is easy to observe that

$$\lambda_n = \frac{c}{r_n^2} \tag{9}$$

with  $c$  defined in Theorem. Define

$$G_n(t) = \int_{B(x_n, r_n)} u(x, t) \phi_n(x - x_n) dx.$$

Let  $\nu_n(x)$  denote the outward unit normal to  $B(0, r_n)$  at  $x \in \partial B(0, r_n)$ . By (2) and the fact that  $\phi_n = 0$  and  $\partial\phi_n/\partial\nu_n \leq 0$  on  $\partial B(0, r_n)$  with the unit normal vector  $\nu_n$ , we obtain

$$G_n(t) \geq G_n(0) + \int_0^t \int_{B(x_n, r_n)} (-\lambda_n u^m(x, s)\phi(x) + u^p(x, s)\phi(x)) dx ds.$$

Put

$$h_n(s) = \begin{cases} -\lambda_n s^m + s^p, & s \geq \left(\frac{m\lambda_n}{p}\right)^{\frac{1}{p-m}}, \\ -\lambda_n \left(\frac{m\lambda_n}{p}\right)^{\frac{m}{p-m}} + \left(\frac{m\lambda_n}{p}\right)^{\frac{p}{p-m}}, & 0 \leq s \leq \left(\frac{m\lambda_n}{p}\right)^{\frac{1}{p-m}}, \end{cases} \quad (10)$$

similarly as in [5]. Since  $h_n$  is convex, we obtain

$$G_n(t) \geq G_n(0) + \int_0^t h_n(G_n(s)) ds. \quad (11)$$

by Jensen's inequality. Let us consider the system of ordinary differential equations

$$\begin{cases} g'_n(t) = h_n(g_n(t)), \\ g_n(0) = G_n(0) \geq b_n. \end{cases} \quad (12)$$

Define  $T_{g_n} = \sup\{t \geq 0 : g_n(t) < \infty\}$  and  $T_{G_n} = \sup\{t \geq 0 : G_n(t) < \infty\}$ . Since  $g_n$  satisfies

$$g_n(t) = g_n(0) + \int_0^t h_n(g_n(s)) ds,$$

and from Lemma, we obtain  $G_n \geq g_n$  and  $T_{g_n} \geq T_{G_n}$ .

Consider the solutions of (1) with the initial data  $b_n$ . The maximal existence times of the solutions denoted by  $T^*(b_n)$  is estimated as

$$T^*(b_n) = \int_{b_n}^{\infty} \frac{d\xi}{\xi^p}.$$

Note that  $\lim_{n \rightarrow \infty} T^*(b_n) = 0$ . From (3) we may assume that there exist  $n_0 \geq 0$  such that

$$\frac{1}{r_n^2 b_n^{p-m}} < \varepsilon$$

for  $n \geq n_0$  and  $\varepsilon \in (0, 1/c)$ . From (9) we see that

$$\lambda_n b_n^m < c\varepsilon b_n^p,$$

and

$$\lambda_n \xi^m < c\varepsilon \xi^p \quad (13)$$

for  $\xi \geq b_n$  and  $n \geq n_0$ . Since  $b_n \geq (m\lambda_n/p)^{1/(p-m)}$  by (13), we have

$$T_{g_n} = \int_{b_n}^{\infty} \frac{d\xi}{h_n(\xi)} = \int_{b_n}^{\infty} \frac{d\xi}{-\lambda_n \xi^m + \xi^p}$$

for  $n \geq n_0$  by (13). Thus we see that

$$\frac{T^*(b_n)}{T_{g_n}} = \frac{\int_{b_n}^{\infty} d\xi/\xi^p}{\int_{b_n}^{\infty} d\xi/(-\lambda_n \xi^m + \xi^p)} > \frac{\int_{b_n}^{\infty} d\xi/\xi^p}{\int_{b_n}^{\infty} d\xi/\{(1-c\varepsilon)\xi^p\}} > 1 - c\varepsilon \quad (14)$$

for  $n \geq n_0$ . Thus we obtain

$$\lim_{m \rightarrow \infty} \frac{T^*(b_n)}{T_{g_n}} \geq 1 - c\varepsilon > 0.$$

Noting that  $\lim_{n \rightarrow \infty} T^*(b_n) = 0$ , we see that  $\lim_{n \rightarrow \infty} T_{g_n} = 0$ . Again we get  $T_{G_n} \rightarrow 0$  as  $n \rightarrow \infty$ . By the definition of the weak solution we have  $T^* = 0$ .  $\square$

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